



<u>Wei Li</u>, Thomas Brumme, Thomas Heine Chair of Theoretical Chemistry, TU Dresden **Relaxation effects in twisted transition metal dichalcogenide heterostructures**

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Why twisted: Moiré





Florian Arnold. (July 5, 2022). Moiré clocks (Video File). Retrieved from https://www.youtube.com/watch?v=uxoKzBPFbrs&list=PL2LyfOO_UvEx5pfK mWlii_hmRBmZwJVG6&index=16.



Relaxation effects

Experimental observation: atomic reconstructions



Twisted bilayer MoS₂^[1]

Bilayer MoS₂ at $\theta = 0.3^{\circ[2]}$



[1] Nat. Comm., **2022**, 13, 3898.
[2] Nat. Nanotechnol., **2020**, 15, 592.



Unrelaxed MoS₂/MoSe₂ at 0°



Fig 1. (a) 3 high-symmetry stackings and transition stacking regions. (b) Interlayer distance and binding energy of corresponding high-symmetry stackings.





Relaxed MoS₂/MoSe₂ at 0°



d (Å)



- Out-of-plane displacement Z_{ML}
- Average Out-of-plane displacement \bar{Z}_{ML}
- Interlayer distance d





Lattice reconstruction and corrugation









Strain energy and vdW energy







Twisted MoS₂/MoSe₂ and MoS₂/WS₂



 $\bigcirc R_h^h \bigvee R_h^M \bigwedge R_h^X$





 \bigcirc H_h^M \bigcirc H_h^h \bigwedge H_h^X



Small vs. large twist angles in MoS₂/WS₂



 Fig 1. (a) Out-of-plane corrugation of Mo along diagonal direction, (b) magnitude of corrugation of in each layer, (c) binding/strain energy at different twist angles.





Spin-Orbit Coupling effect







Summary

- Significant lattice reconstruction
 - Domain formation



Strain energy cost and van der Waals energy gain





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How: Multiscale approach



- Geometry optimization performed by Force-Field^{[1][2]}
- Electronic properties calculated by DFTB^[3]
- Force-Field validated by DFT

- Lattice size: 10¹ 10³ Å
- Number of atoms: 50 5×10⁵



[1] J. Appl. Phys. 2013, 114, 064307.

[2] J. Phys. Chem. C **2019**, 123, 9770.

[3] J. Chem. Theory Comput. **2022**, 18, 4472.



Generating models

$$\begin{bmatrix} \tilde{a}_{1}^{T} \\ \tilde{a}_{2}^{T} \end{bmatrix} = M_{a} \begin{bmatrix} a_{1}^{T} \\ a_{2}^{T} \end{bmatrix}, \begin{bmatrix} \tilde{b}_{1}^{T} \\ \tilde{b}_{2}^{T} \end{bmatrix} = RM_{b} \begin{bmatrix} b_{1}^{T} \\ b_{2}^{T} \end{bmatrix}$$

$$\blacksquare \text{ Hexagonal symmetry}$$

$$M_{j} = \begin{bmatrix} p_{j} & q_{j} \\ -q_{j} & p_{j} - q_{j} \end{bmatrix}$$

$$\blacksquare \text{ Coincidence}$$

$$|\tilde{a}_{i} - \tilde{b}_{i}| < tolerance$$





Geometry optimization: Force-field method







Computational details

- Geometry optimisation:
 - Conjugate gradient method in LAMMPS
- Single point calculations (DFT):
 - GGA-PBE/TS in FHI-aims
- Single point calculations (DFTB):
 - SCC-DFTB theory level in DFTB+
 - QUASINANO2013 Slater-Koster parameters
 - Monkhorst Pack k-space





Boundary condition: flake







Unrelaxed MoS₂/MoSe₂ at 60°



Fig 1. (a) 3 high-symmetry stackings and transition stacking regions. (b) Interlayer distance and binging energy of corresponding high-symmetry stackings.





Backup



 Fig 1. Interlayer distance and binding energy of high-symmetry stackings of MoS₂/WS₂



