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Spin Orbit coupling implementation in DFTB/GFN-xTB

Kohn-Sham equations $E = E[\rho(\mathbf{r})]$





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Expanding $E_{
ho}$ at ho_0 to second (or third) order in fluctuation $\delta
ho$





Kohn-Sham equations $E = E[\rho(\mathbf{r})] \rightarrow \rho = \rho_0 + \delta \rho$

Expanding E_{ρ} at ρ_0 to second (or third) order in fluctuation $\delta \rho$

•
$$E^0[\rho_0] = -\frac{1}{2} \int \int \frac{\rho_0(\mathbf{r})\rho_0(\mathbf{r}')}{|\mathbf{r}-\mathbf{r}'|} d\mathbf{r} d\mathbf{r}' - \int V^{xc}[\rho_0]\rho_0(\mathbf{r}) d\mathbf{r} + E_{\mathsf{xc}}[\rho_0] + \frac{1}{2} \sum_{A \neq B} \frac{Z_A Z_B}{R_{AB}}$$

•
$$E_1[\rho_0, \delta\rho] = \sum_i^{\text{occ}} f_i \langle \psi_i | -\frac{1}{2} \nabla^2 + V_{\text{ext}} + \int \frac{\rho_0(\mathbf{r})}{|\mathbf{r} - \mathbf{r}'|} + V_{xc}[\rho_0] |\psi_i \rangle$$

•
$$E^2[\rho_0, \delta \rho^2] = \frac{1}{2} \int \int \left(\frac{1}{|\mathbf{r} - \mathbf{r}'|} + \frac{\delta^2 E_{\mathsf{xc}}}{\delta \rho(\mathbf{r}) \delta \rho(\mathbf{r}')} \right) \delta \rho(\mathbf{r}) \delta \rho(\mathbf{r}') \, d\mathbf{r} \, d\mathbf{r}'$$





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Expanding $E_{
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•
$$E^2[\rho_0, \delta \rho^2] = \frac{1}{2} \int \int \left(\frac{1}{|\mathbf{r} - \mathbf{r}'|} + \frac{\delta^2 E_{\mathsf{xc}}}{\delta \rho(\mathbf{r}) \delta \rho(\mathbf{r}')} \right) \delta \rho(\mathbf{r}) \delta \rho(\mathbf{r}') \, d\mathbf{r} \, d\mathbf{r}'$$



Non-SCC (?) DFTB (DFTB1)

Expanding $E_{
ho}$ at ho_0 to second (or third) order in fluctuation $\delta
ho$

$$E_{\text{tot}}[\rho_0 + \delta\rho] = E^0[\rho_0] + E^1[\rho_0, \delta\rho] + E^2[\rho_0, \delta\rho^2] + O(\delta\rho^3)$$
$$= \sum_{I < J} V_{IJ}^{\text{rep}}(\mathbf{r}_{IJ}) + \sum_i^{\text{occ}} f_i \langle \psi_i | \hat{H}[\rho_0] | \psi_i \rangle + \dots$$





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- minimal "pseudo" atomic basis set $\psi_i = \sum_{\mu} c_{\mu i} \varphi_{\mu}$
- reference density ρ_0
- tabulated as function of distance





Non-SCC (?) DFTB (DFTB1)

$$E_{\text{Non-SCC}} = \sum_{I < J} V_{IJ}^{\text{rep}}(\mathbf{r}_{IJ}) + \sum_{i}^{\text{occ}} f_i \langle \psi_i | \hat{H}[\rho_0] | \psi_i \rangle$$

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Expanding $E_{
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$$E^2[\rho_0, \delta \rho^2] = \frac{1}{2} \int \int \left(\frac{1}{|\mathbf{r} - \mathbf{r}'|} + \frac{\delta^2 E_{\mathsf{XC}}}{\delta \rho(\mathbf{r}) \delta \rho(\mathbf{r}')} \right) \delta \rho(\mathbf{r}) \delta \rho(\mathbf{r}') \, d\mathbf{r} \, d\mathbf{r}'$$





Expanding $E_{
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 $E_{\text{tot}}[\rho_0 + \delta\rho] = E^0[\rho_0] + E^1[\rho_0, \delta\rho] + E^2[\rho_0, \delta\rho^2] + O(\delta\rho^3)$

•
$$E^{2}[\rho_{0}, \delta\rho^{2}] = \frac{1}{2} \int \int \left(\frac{1}{|\mathbf{r}-\mathbf{r}'|} + \frac{\delta^{2} E_{\mathrm{xc}}}{\delta\rho(\mathbf{r})\delta\rho(\mathbf{r}')} \right) \delta\rho(\mathbf{r}) \delta\rho(\mathbf{r}') d\mathbf{r} d\mathbf{r}'$$

$$\int \delta\rho(\mathbf{r}) = \sum_{I} \Delta q_{I} \delta\rho_{I}(\mathbf{r})$$

$$E^{2}[\rho_{0}, \delta\rho^{2}] = \frac{1}{2} \sum_{I} \left(\sum_{i=1}^{N} \Delta q_{I} \delta\rho_{I}(\mathbf{r}) \right) \delta\rho(\mathbf{r}') d\mathbf{r} d\mathbf{r}'$$

• $E^2[\rho_0, \delta \rho^2] \approx \frac{1}{2} \sum_{IJ} \gamma_{IJ}(R_{IJ}) \Delta q_I \Delta q_J$







- $E_{\text{DFTB2(3)}} = E_{\text{Non-SCC}} + \frac{1}{2} \sum_{IJ} \gamma_{IJ} (R_{IJ}) \Delta q_I \Delta q_J + \frac{1}{3} \sum_{IJ} \Delta q_I^2 \Delta q_J \Gamma_{IJ}$
- $H_{\mu\nu} = H^0_{\mu\nu} + H^2_{\mu\nu}[\gamma^h, \Delta q] + H^3_{\mu\nu}[\Gamma, \Delta q], \quad \mu \in I, \nu \in J$
- $q_I = \sum_i f_i \int_{V_I} |\psi_i(\mathbf{r})|^2 d^3 r = \frac{1}{2} \sum_i^{\text{occ}} f_i \sum_{\mu \in I} \sum_{\nu} \left(c^*_{\mu i} c_{\nu i} S_{\mu\nu} + c^*_{\nu i} c_{\mu i} S_{\nu\mu} \right)$



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Charges iterated until self consistency has been reached



Self-consistent charge (SCC) iteration



Spin-Orbit Coupling (SOC)



$$\Delta E_{SO} \qquad \Delta E_{SO} = \xi (L \cdot S)$$

- > SOC hamiltonian in Dirac Equation $\hat{H}^{SOC} = -\frac{e\hbar}{4m^2c^2}\mathbf{\sigma}\cdot[E_f\times\hat{p}]$
- > Spherical potential & static case

$$\hat{H}^{soc} = -\frac{\hbar}{4m^2c^2} \frac{1}{r} \frac{dV}{dr} \boldsymbol{\sigma} \cdot [\hat{r} \times \hat{p}]$$

 $\hat{H}^{soc} = \xi(\hat{L}\cdot\hat{S})$



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 $\hat{H}^{soc} = -\frac{\hbar}{4m^2c^2}\frac{1}{r}\frac{dV}{dr}\boldsymbol{\sigma}\cdot[\hat{r}\times\hat{p}]$

 $\hat{H}^{soc} = \xi(\hat{L} \cdot \hat{S})$

See J. Chem. Theory Comput. 18, 4472 (2022)



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• Single-particle on-site spin-orbit interaction

$$\begin{split} \hat{H}_{SO} &= \frac{\zeta}{2} \mathbf{L} \cdot \mathbf{S} = \frac{\zeta}{2} (\hat{L}_x \sigma_x + \hat{L}_y \sigma_y + \hat{L}_z \sigma_z) \\ &= \frac{\zeta}{2} \begin{pmatrix} \hat{L}_z & \hat{L}_x - i\hat{L}_y \\ \hat{L}_x + i\hat{L}_y & \hat{L}_z \end{pmatrix} = \frac{\zeta}{2} \begin{pmatrix} \hat{L}_z & \hat{L}_- \\ \hat{L}_+ & \hat{L}_z \end{pmatrix} \end{split}$$

Pauli matrices:







Single-particle on-site spin-orbit interaction

$$\hat{H}_{SO} = \frac{\zeta}{2} \mathbf{L} \cdot \mathbf{S} = \frac{\zeta}{2} (\hat{L}_x \sigma_x + \hat{L}_y \sigma_y + \hat{L}_z \sigma_z) \qquad \sigma_y = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}$$
$$= \frac{\zeta}{2} \begin{pmatrix} \hat{L}_z & \hat{L}_x - i\hat{L}_y \\ \hat{L}_x + i\hat{L}_y & \hat{L}_z \end{pmatrix} = \frac{\zeta}{2} \begin{pmatrix} \hat{L}_z & \hat{L}_- \\ \hat{L}_+ & \hat{L}_z \end{pmatrix} \qquad \sigma_z = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$$

Pauli matrices:

$$\sigma_x = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$$

$$= \begin{pmatrix} 0 & -i \end{pmatrix}$$

- $\hat{L}_{z}|Y_{l}^{\pm m}\rangle = \pm ml|Y_{l}^{\pm m}\rangle$, $\hat{L}_{\pm}|Y_{l}^{\pm m}\rangle = \sqrt{l(l+1) - m(m\pm 1)}|Y_{l}^{\pm m}\rangle$ - $\varphi_{\mu}(\mathbf{r} - \mathbf{R}_{I}) = R_{\mu}(r)\widetilde{Y}_{\mu}(\theta,\varphi)(\mu \in I)$, where $\widetilde{Y} \propto Y_{lm} \pm Y_{lm}^{*}$





• 2-component spinor wavefunctions

$$\psi_i = \sum_{\mu} \begin{pmatrix} c^{\alpha}_{\mu i} \\ c^{\beta}_{\mu i} \end{pmatrix} \varphi_{\mu}$$

Non-collinear Hamiltonian

$$\hat{H} = \left(\hat{H}^{0}_{\mu\nu} + \hat{H}^{2}_{\mu\nu} + \hat{H}^{3}_{\mu\nu}\right) \otimes \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} + \begin{bmatrix} \xi_{I}^{l} \begin{pmatrix} \hat{L}_{z} & \hat{L}_{-} \\ \hat{L}_{+} & -\hat{L}_{z} \end{bmatrix}_{l} + \xi_{J}^{l'} \begin{pmatrix} \hat{L}_{z} & \hat{L}_{-} \\ \hat{L}_{+} & -\hat{L}_{z} \end{bmatrix}_{l'} \end{bmatrix}$$



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• Secular equation

$$\sum_{\mu} \begin{pmatrix} H^{\alpha\alpha}_{\mu\nu} - \epsilon_i S_{\mu\nu} & H^{\alpha\beta}_{\mu\nu} \\ H^{\beta\alpha}_{\mu\nu} & H^{\beta\beta}_{\mu\nu} - \epsilon_i S_{\mu\nu} \end{pmatrix} \begin{pmatrix} c^{\alpha}_{\mu i} \\ c^{\beta}_{\mu i} \end{pmatrix} = 0, \quad H^{\sigma\sigma'}_{\mu\nu} = \langle \varphi_{\mu} | \hat{H}^{\sigma\sigma'} | \varphi_{\nu} \rangle$$



• Density matrix

$$\rho(\mathbf{r}) = \sum_{i}^{\mathsf{occ}} f_i \begin{pmatrix} \psi_i^{\alpha*} \\ \psi_i^{\beta*} \end{pmatrix} \begin{pmatrix} \psi_i^{\alpha} \\ \psi_i^{\beta} \end{pmatrix} = \begin{pmatrix} \rho^{\alpha\alpha} & \rho^{\alpha\beta} \\ \rho^{\beta\alpha} & \rho^{\beta\beta} \end{pmatrix}$$
$$= n(\mathbf{r}) \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} + m^x(\mathbf{r})\sigma_x + m^y(\mathbf{r})\sigma_y + m^z(\mathbf{r})\sigma_z$$



• Electron and magnetization densities

$$n(\mathbf{r}) = \frac{1}{2} \operatorname{Re} \left(\rho^{\alpha \alpha} + \rho^{\beta \beta} \right), \quad m^{z}(\mathbf{r}) = \frac{1}{2} \operatorname{Re} \left(\rho^{\alpha \alpha} - \rho^{\beta \beta} \right)$$





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$$= n(\mathbf{r}) \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} + m^x(\mathbf{r})\sigma_x + m^y(\mathbf{r})\sigma_y + m^z(\mathbf{r})\sigma_z$$



• Electron and magnetization densities

$$n(\mathbf{r}) = \frac{1}{2} \operatorname{Re} \left(\rho^{\alpha \alpha} + \rho^{\beta \beta} \right), \quad m^{z}(\mathbf{r}) = \frac{1}{2} \operatorname{Re} \left(\rho^{\alpha \alpha} - \rho^{\beta \beta} \right)$$

• Mulliken charges

$$q_{\mu} = \frac{1}{2} \left(q_{\mu}^{\alpha \alpha} + q_{\mu}^{\beta \beta} \right), \quad q_{\mu}^{\sigma \sigma'} = \sum_{\nu} \rho^{\sigma \sigma'} S_{\mu \nu}$$



• Energy: $E_{SO} = Tr(\rho H_{SO})$

$$E_{\mathsf{SO}} = \sum_{\mu\nu} \left[H^{\alpha\alpha}_{\mathsf{SO}} \rho^{\alpha\alpha}_{\nu\mu} + H^{\beta\beta}_{\mathsf{SO}} \rho^{\beta\beta}_{\nu\mu} + 2 \operatorname{Re} \left(H^{\alpha\beta}_{\mathsf{SO}} \rho^{\alpha\beta}_{\nu\mu} \right) \right]$$

• Forces : $\mathbf{F}_{\mathsf{SO},I} = -\sum_{\sigma\sigma'} \frac{\partial E^{\sigma\sigma'}_{\mathsf{SO}}}{\partial \mathbf{R}_{I}}$

$$\mathbf{F}_{\mathsf{SO},I} = \sum_{\mu\nu,\mathbf{R}} \left[\sum_{\sigma\sigma'} \rho_{\mu\nu}^{\sigma\sigma'}(\mathbf{R}) \frac{\partial H_{\mathsf{SO}}^{\sigma\sigma'}(\mathbf{R})}{\partial \mathbf{R}_{I}} - \sum_{\sigma\sigma'} \rho_{\mathsf{SO}}^{\epsilon,\sigma\sigma'}(\mathbf{R}) \frac{\partial S_{\mu\nu}(\mathbf{R})}{\partial \mathbf{R}_{I}} \right]$$



m / **T**

• Energy:
$$E_{SO} = \operatorname{Tr}(\rho H_{SO})$$

 $E_{SO} = \sum_{\mu\nu} \left[H_{SO}^{\alpha\alpha} \rho_{\nu\mu}^{\alpha\alpha} + H_{SO}^{\beta\beta} \rho_{\nu\mu}^{\beta\beta} + 2\operatorname{Re}\left(H_{SO}^{\alpha\beta} \rho_{\nu\mu}^{\alpha\beta}\right) \right]$
• Forces: $\mathbf{F}_{SO,I} = -\sum_{\sigma\sigma'} \frac{\partial E_{SO}^{\sigma\sigma'}}{\partial \mathbf{R}_{I}}$
 $\mathbf{F}_{SO,I} = \sum_{\mu\nu,\mathbf{R}} \left[\sum_{\sigma\sigma'} \rho_{\mu\nu}^{\sigma\sigma'}(\mathbf{R}) \underbrace{\frac{\partial H_{SO}^{\sigma\sigma'}(\mathbf{R})}{\partial \mathbf{R}_{I}}}_{0} - \underbrace{\sum_{\sigma\sigma'} \rho_{SO}^{\epsilon,\sigma\sigma'}(\mathbf{R}) \frac{\partial S_{\mu\nu}(\mathbf{R})}{\partial \mathbf{R}_{I}}}_{0} \right]$
- No explicit contribution $\frac{\mathbf{Bi}_{2} \quad \mathbf{Bond \, Length}\left(\mathbf{A}\right) \quad \mathbf{Frequency}\left(\mathbf{cm}^{-1}\right)}{\mathbf{W} \operatorname{SOC}} \quad 1.98 \quad 1366 \\ \mathbf{W}/t \operatorname{SOC} \quad 2.02 \quad 1336 \end{array}$

Table: calculated with GFN1-xTB in AMS/DFTB



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• Periodic boundary conditions

$$\phi_{\mu}^{\mathbf{k}}(\mathbf{r}) = \frac{1}{\sqrt{N}} \sum_{\mathbf{R}} e^{i\mathbf{k}\cdot\mathbf{R}} \varphi_{\mu}(\mathbf{r} - \mathbf{R})$$



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• Fatbands

$$n_{i\mathbf{k}}^{\mu} = \sum_{\nu} \left| \langle \psi_{\mu}^{\mathbf{k}} | \phi_{\nu}^{\mathbf{k}} \rangle \right|^{2} = \frac{1}{2} \sum_{\nu} \operatorname{Re} \left(c_{\mu i \mathbf{k}}^{\alpha *} c_{\nu i \mathbf{k}}^{\alpha} S_{\mu \nu} + c_{\mu i \mathbf{k}}^{\beta *} c_{\nu i \mathbf{k}}^{\beta} S_{\mu \nu} \right)$$

• Spin texture $S_{i\mathbf{k}} = \langle \psi_{i\mathbf{k}} | \sigma | \psi_{i\mathbf{k}} \rangle$

$$m_{i\mathbf{k}}^{\mu,z} = \frac{1}{2} \sum_{\nu} \operatorname{Re} \left(c_{\mu i\mathbf{k}}^{\alpha*} c_{\nu i\mathbf{k}}^{\alpha} S_{\mu\nu} - c_{\mu i\mathbf{k}}^{\beta*} c_{\nu i\mathbf{k}}^{\beta} S_{\mu\nu} \right)$$



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Benchmark calculations



- Me functionalized Bi(111) topological insulators
- QUASINANO2013 Slater-Koster parameters
- Visualized in amsbands



Benchmark calculations



- 2D WS₂ with SOC
- Plotted with python
- Spinor Visualization in amsbands to be done



Benchmark calculations

- *H*_{SO} is **transferable** among different parameter sets and methods
- Successful benchmarked on close-shell **molecules** and **materials** such as III-V 3D semiconductors, TMDC 2D crystals, topological insulators, with comparison to DFTB+
- **Regression test** on single point calculation, geometry optimization, and frequency calculation



Scalability

TIH chain







THANK YOU





Backup: GFN-xTB

 $\psi_i = \sum_{\mu} c_{\mu i} \varphi_{\mu}(\zeta, \mathsf{STO} - mG) \mathsf{b}$

$$\begin{split} H_{\mu\nu} &= K_{IJ} \frac{1}{2} (k_l + k_{l'}) \frac{1}{2} (h_J + h_{J'}) S_{\mu\nu} (1 + k_{EN} \Delta E_N^2) \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} + \Pi(R_{I,l'}) \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} \\ &+ \frac{1}{2} S_{\mu\nu} \sum_c \sum_{l''} \left(Y_{IC,l''} + Y_{JC,l''} \right) P_{l''}^c \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} + \frac{1}{2} S_{\mu\nu} (q_I^2 + q_J^2) \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} \\ &+ \frac{1}{2} S_{\mu\nu} \left[\epsilon_\mu \begin{pmatrix} L_z & L_- \\ L_+ & -L_z \end{pmatrix} + \epsilon_v \begin{pmatrix} L_z & L_- \\ L_+ & -L_z \end{pmatrix} \right] \end{split}$$

